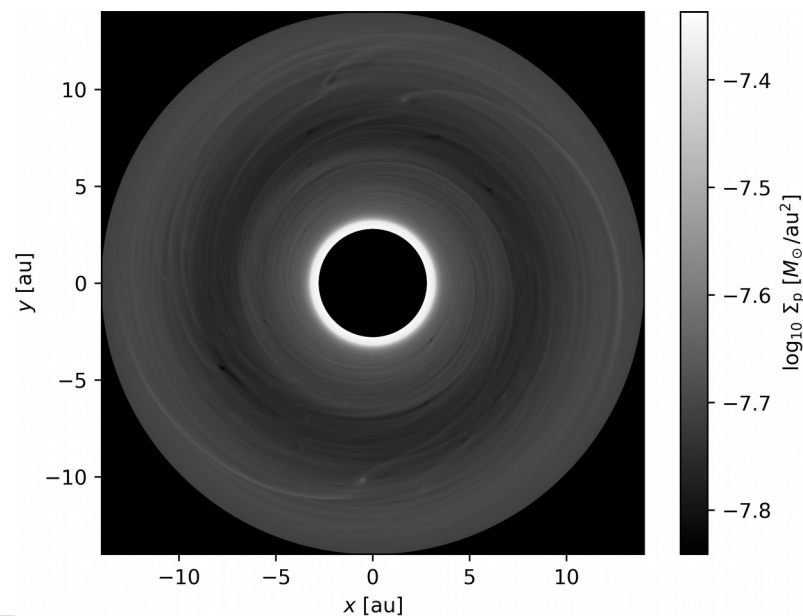
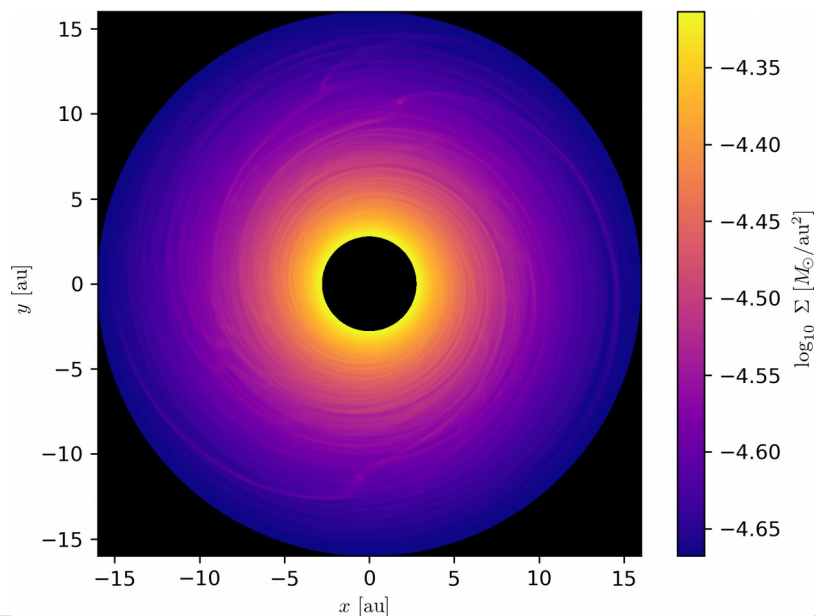


Dynamics of protoplanets embedded in gas/pebble disks

↑ and its dependence on Σ and ν parameters

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**Q: What is the (fastest) mechanism of
giant-planet core formation?**

gas

$$\frac{\partial \Sigma}{\partial t} + \overbrace{\mathbf{v} \cdot \nabla \Sigma}^{\text{convection}} = - \overbrace{\Sigma \nabla \cdot \mathbf{v}}^{\text{expansion}} - \overbrace{\left(\frac{\partial \Sigma}{\partial t} \right)_{\text{acc}}}^{\text{gas accretion}}$$

Masset (2000)
Rein & Spiegel (2015)
Chrenko et al. (2017)

$$\frac{\partial \mathbf{v}}{\partial t} + \mathbf{v} \cdot \nabla \mathbf{v} = - \frac{1}{\Sigma} \nabla P + \overbrace{\frac{1}{\Sigma} \nabla \cdot \mathbb{T}}^{\text{bulk viscosity}} - \overbrace{\frac{\int \rho \nabla \phi dz}{\Sigma}}^{\text{gravity}} + \overbrace{\frac{\Sigma_p \Omega_K}{\Sigma \tau} (\mathbf{u} - \mathbf{v})}^{\text{back-reaction}}$$

$$\frac{\partial E}{\partial t} + \mathbf{v} \cdot \nabla E = -E \nabla \cdot \mathbf{v} - P \nabla \cdot \mathbf{v} + \overbrace{Q_{\text{visc}}}^{\text{dissip.}} + \overbrace{\frac{2\sigma T_{\text{irr}}^4}{\tau_{\text{eff}}}}^{\text{irradiation}} - \overbrace{\frac{2\sigma T^4}{\tau_{\text{eff}}}}^{\text{vertical cooling}} + \overbrace{2H \nabla \cdot \frac{16\sigma \lambda_{\text{lim}} T^3 \nabla T}{\rho_0 \kappa}}^{\text{radiative diffusion}} + \overbrace{\frac{\sum_i GM_i \dot{M}_i}{R_i S_{\text{cell}}}}^{\text{accretion}} \delta(\mathbf{r}_i)$$

$$P = \Sigma \frac{RT}{\mu} = (\gamma - 1)E$$



pebbles

$$\frac{\partial \Sigma_p}{\partial t} + \mathbf{u} \cdot \nabla \Sigma_p = - \Sigma_p \nabla \cdot \mathbf{u} - \overbrace{\left(\frac{\partial \Sigma_p}{\partial t} \right)_{\text{acc}}}^{\text{pebble accretion}}$$

$$\frac{\partial \mathbf{u}}{\partial t} + \mathbf{u} \cdot \nabla \mathbf{u} = - \frac{\int \rho_p \nabla \phi dz}{\Sigma_p} - \overbrace{\frac{\Omega_K}{\tau} (\mathbf{u} - \mathbf{v})}^{\text{aerodynamic drag}}$$

protoplanets

$$\ddot{\mathbf{r}}_i = - \frac{GM_\star}{r_i^3} \mathbf{r}_i - \sum_{j \neq i} \frac{GM_j}{|\mathbf{r}_i - \mathbf{r}_j|^3} (\mathbf{r}_i - \mathbf{r}_j) - \iint \frac{G\Sigma}{|\mathbf{r}_i - \mathbf{r}_{\text{cell}}|^3} (\mathbf{r}_i - \mathbf{r}_{\text{cell}}) r d\theta dr + \overbrace{f_z \hat{z}}^{\text{vertical damping}} \quad \text{for } \forall i$$

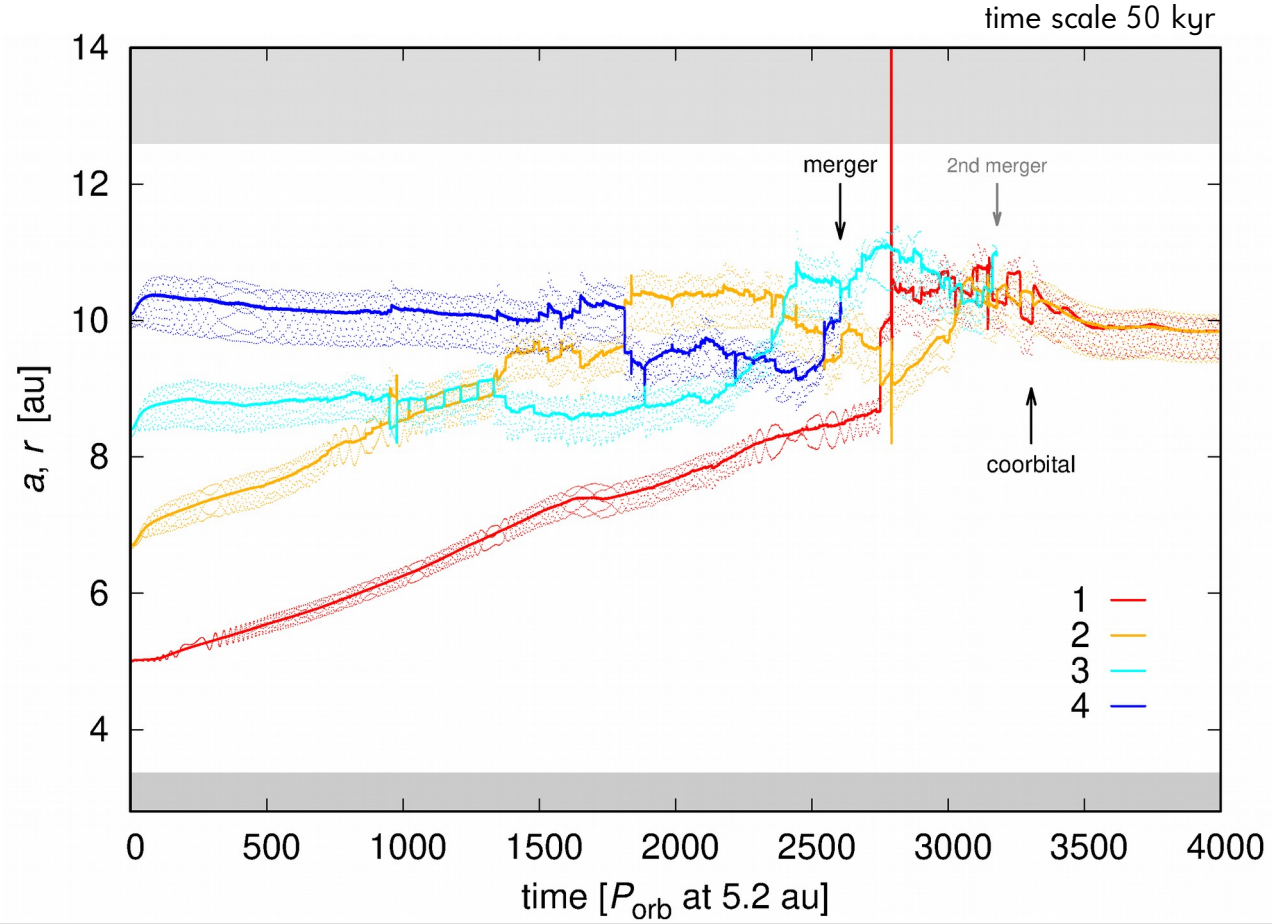
Caselli_nominal

← from Chrenko et al. (2017)

- semimajor axis a vs time t
- convergence zone @ 10 au
- oscillations of r (a.k.a. e)
← **hot-trail** effect
→ no resonant captures
- 2 mergers, 1 coorbital

- BUT see details...

surface density $\Sigma_0 = 750 \text{ g/cm}^2$, slope -0.5,
kinematic viscosity $\nu = 10^{-5}$ [c.u.],
embryo mass $M_{\text{em}} = 3 M_E$,
pebble flux $dM_p/dt = 10^{-4} M_E/\text{yr}$, ...



Detail: Merger

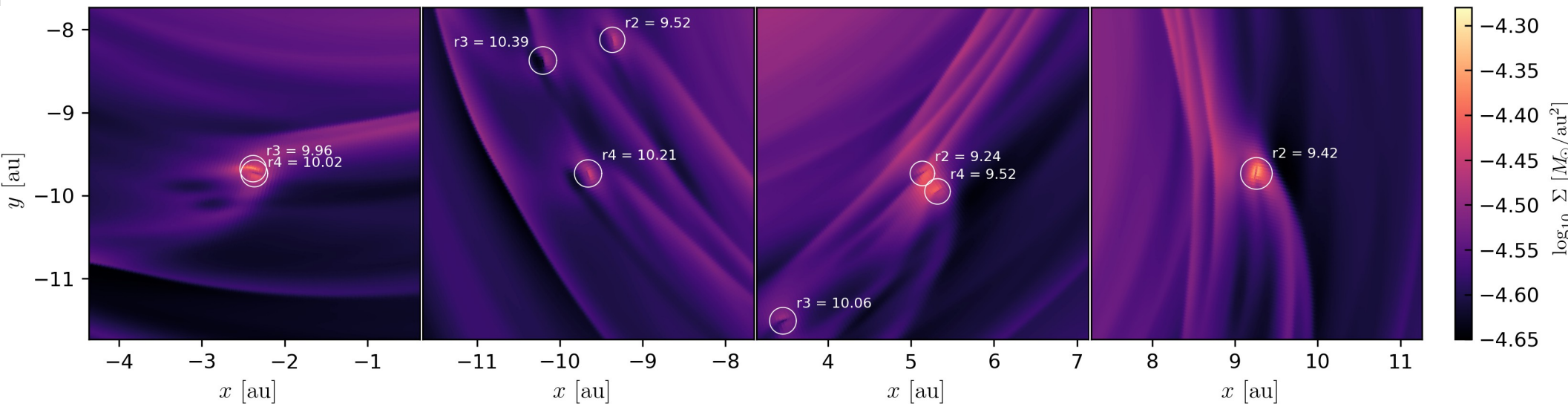
- restarted from sparse output \rightarrow detailed output ($1/20 P_{\text{orb}}$)
- **3-body encounters** are needed for mergers! \leftarrow small statistics

deep encounter...

3 protoplanets...

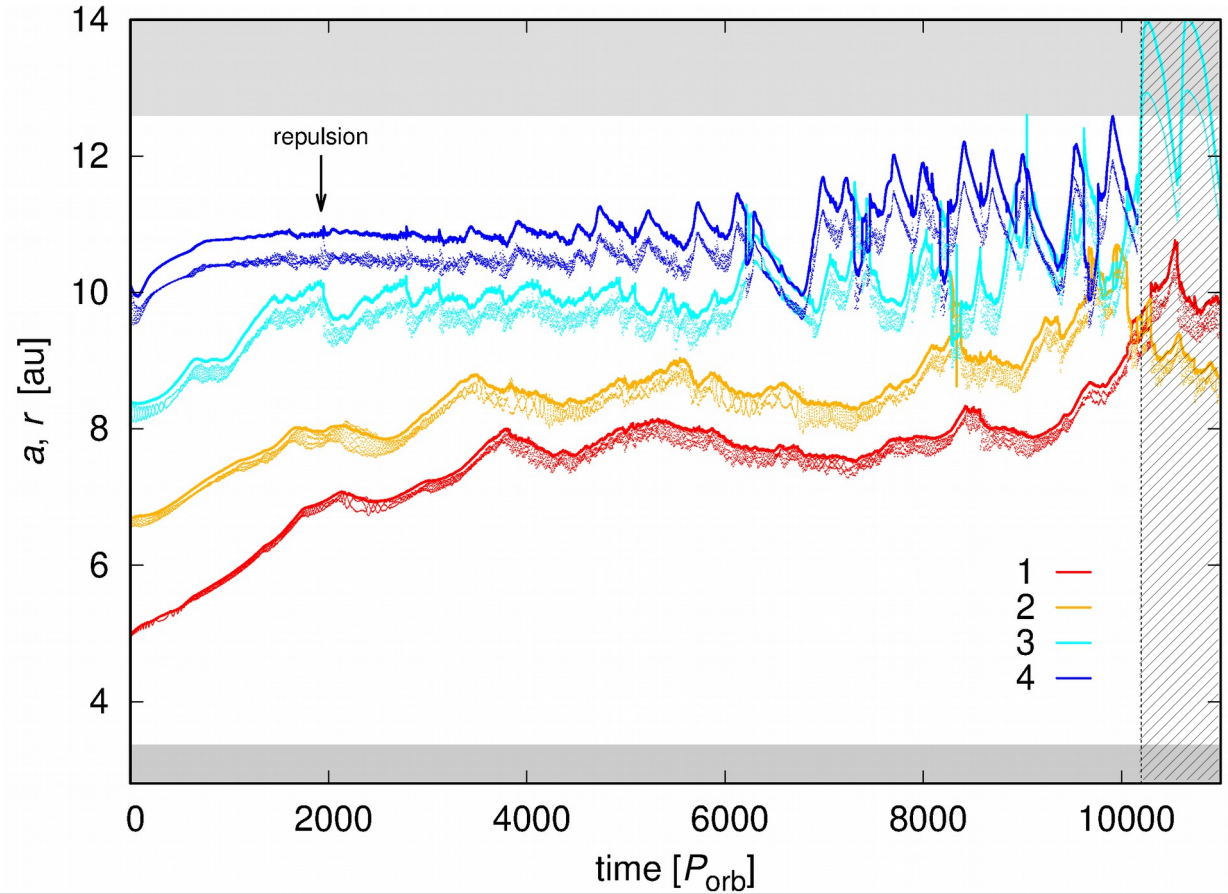
different encounter...

merger



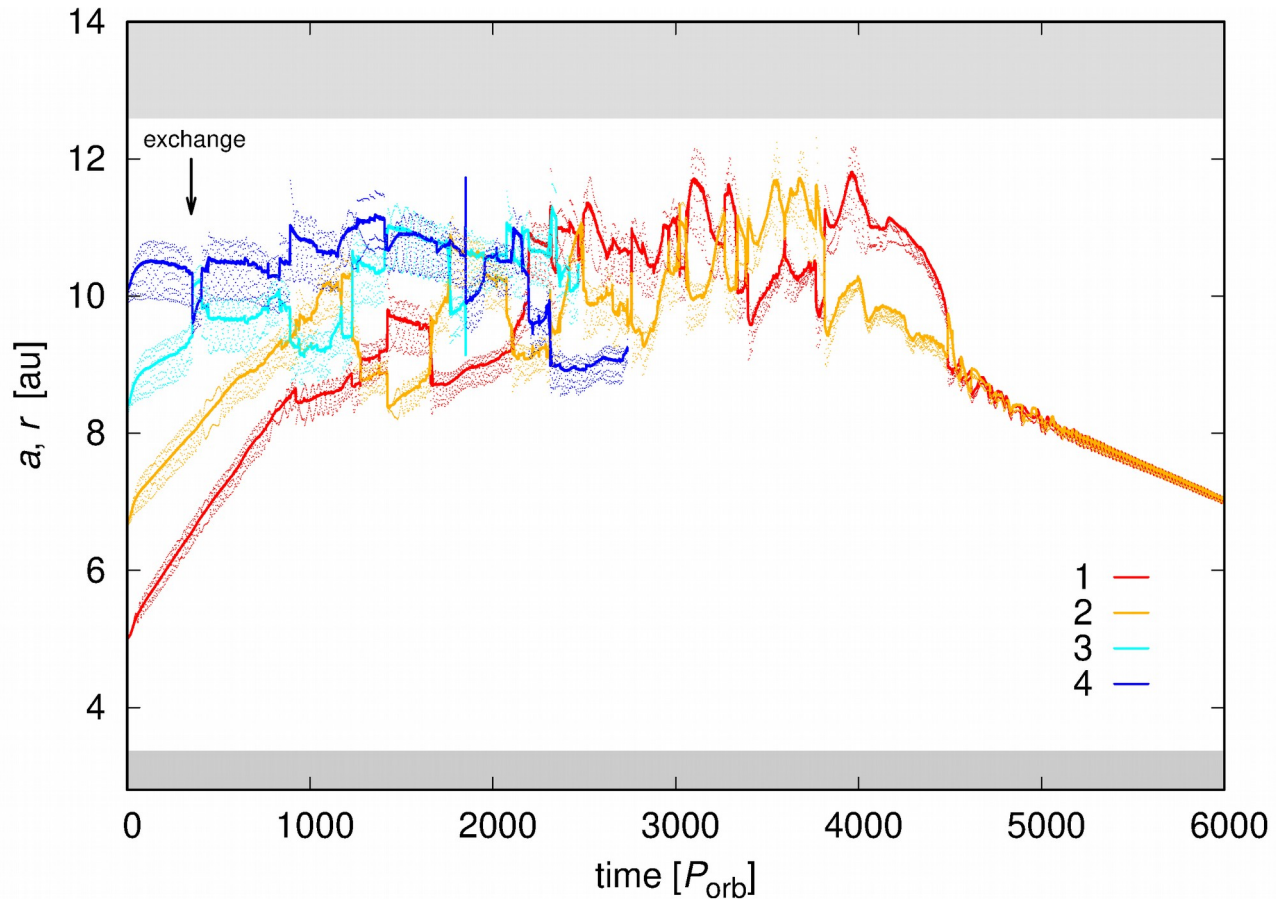
Sigma_3times

- i.e. $\Sigma_0 = 2250 \text{ g/cm}^2$
- substantial interior disk mass
→ a offset from r
- many 'repulsions' (3 and 4)
- HD also determines encounter geometry



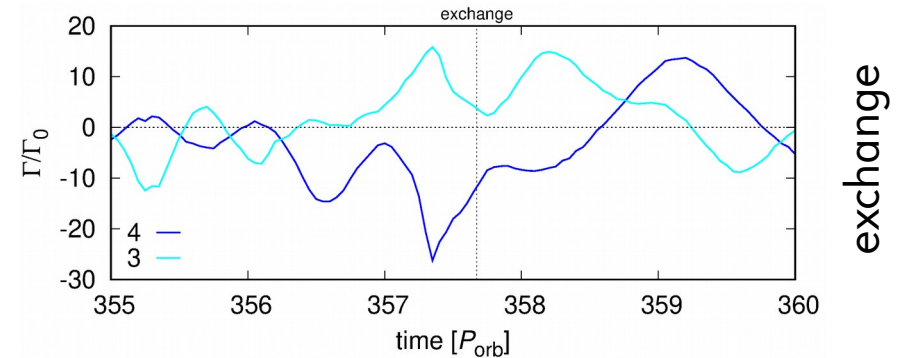
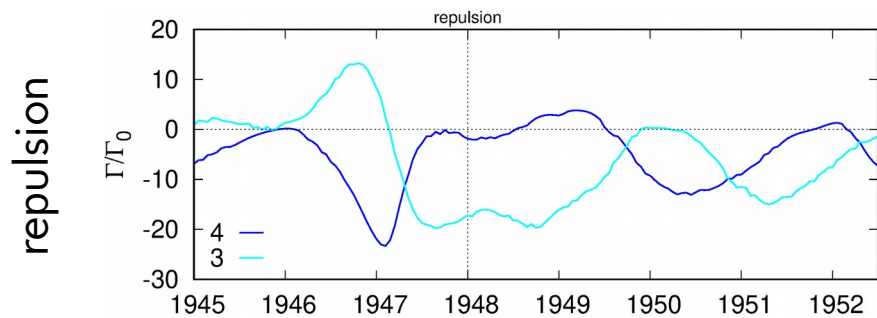
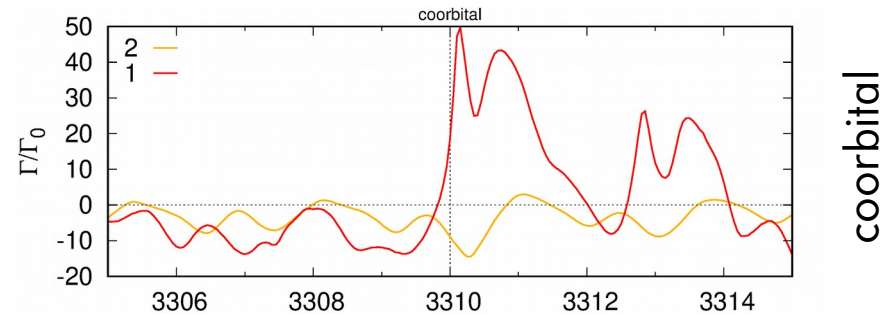
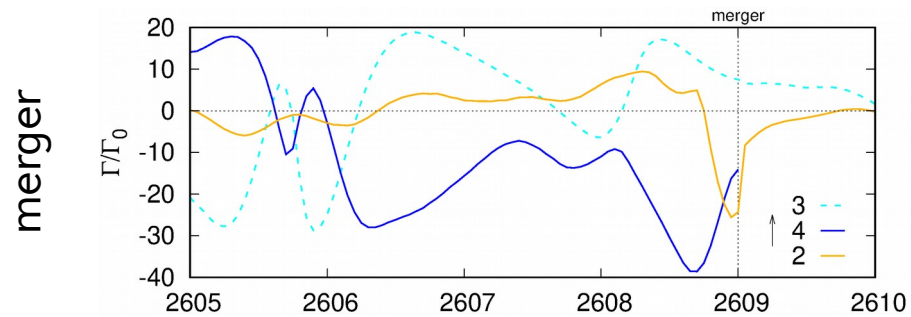
viscosity_1e-6

- i.e. $\nu = 10^{-6}$ [c.u.]
- massive coorbital formation & its **stabilisation**
- too-many-coorbitals problem...
- **pebble isolation** beyond
 - no heating
 - inward migration



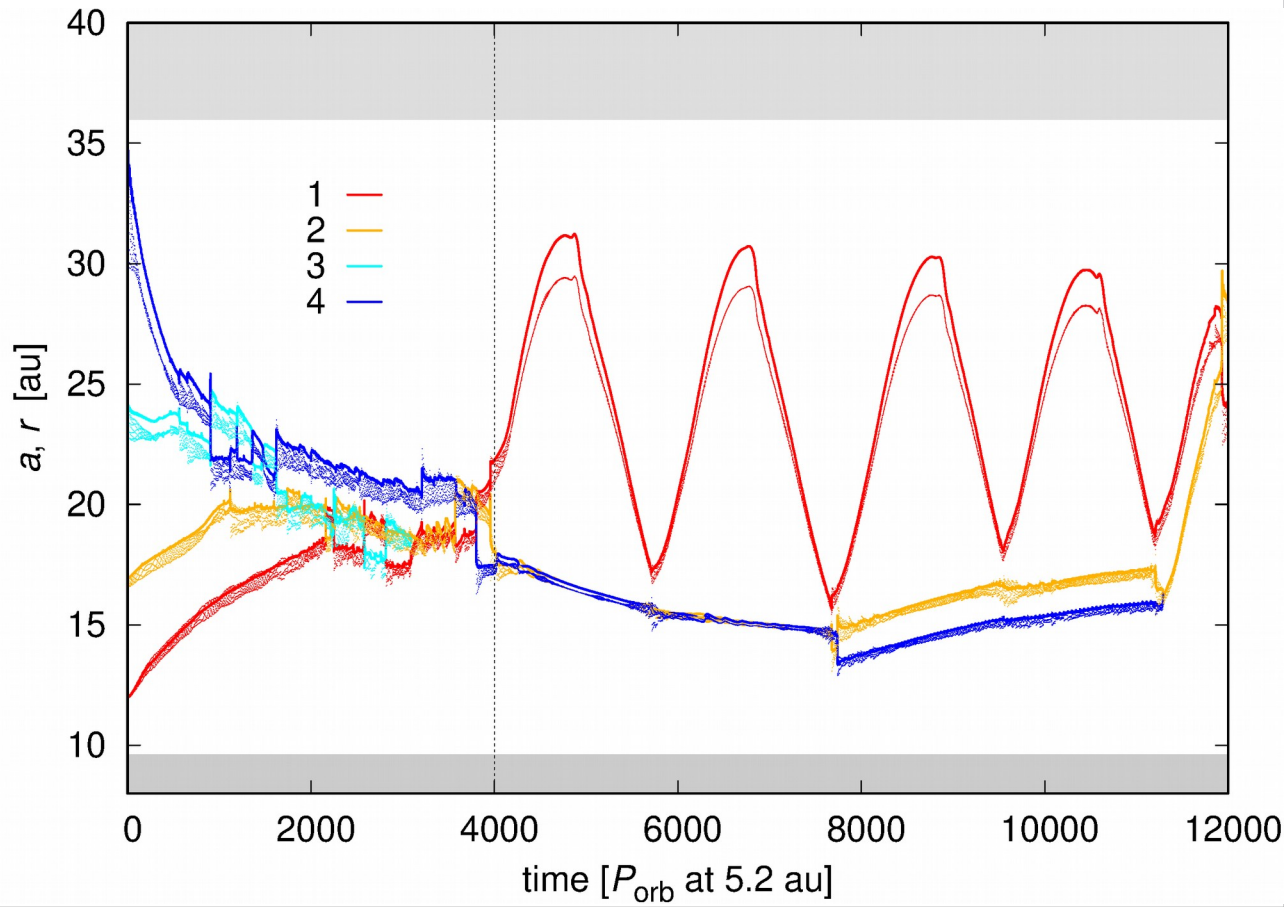
Torque details of 4 events

- normalized disk torque Γ/Γ_0 vs time t , with $\Gamma_0 = (q/h)^2 \Sigma r^4 \Omega^2$
- encounters perturbed by torques, which also determine (next) geometry



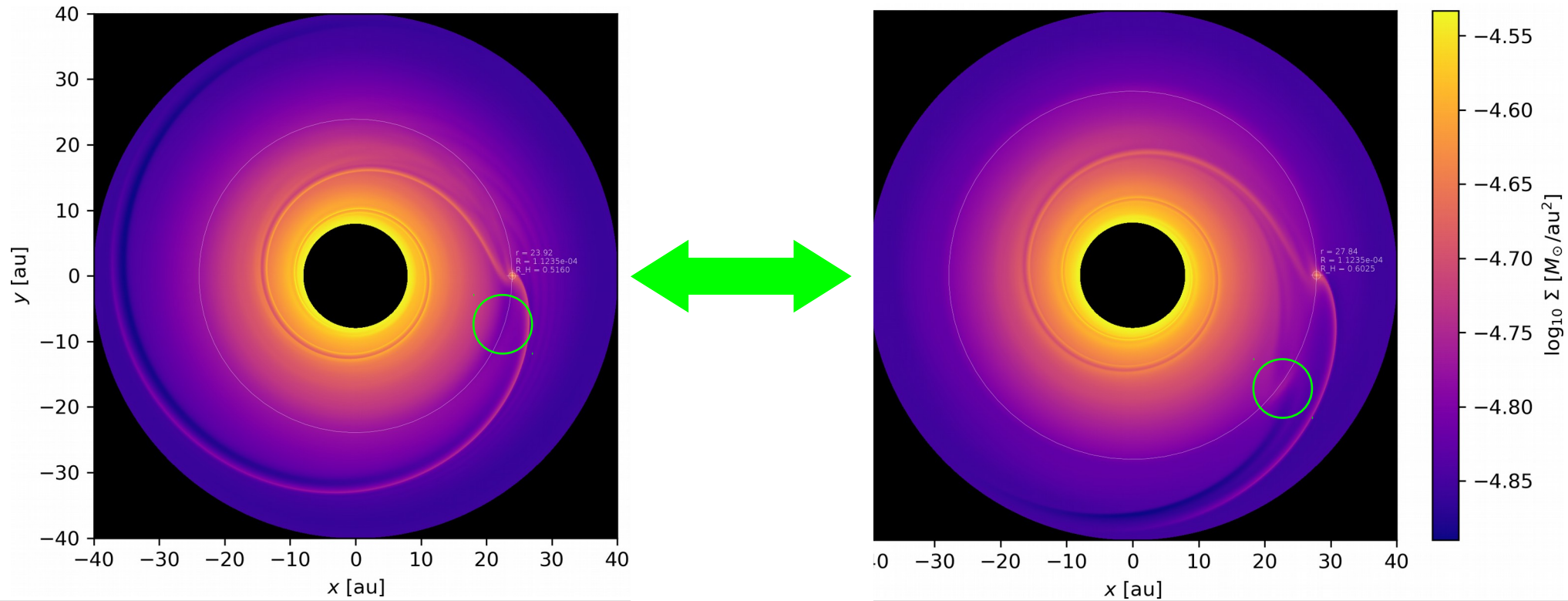
totmass_20ME

- i.e. $M_{\text{em}} = 5 M_{\text{E}}$
- convergence zone @ 20 au
- 3-body interactions again
- coorbital + **embryo 1**
- initial 'kick' → **dynamical torque** drives it to 30 au
- returns back to 20 au & interacts with coorbital
- a solution of too-many-coorbitals problem?



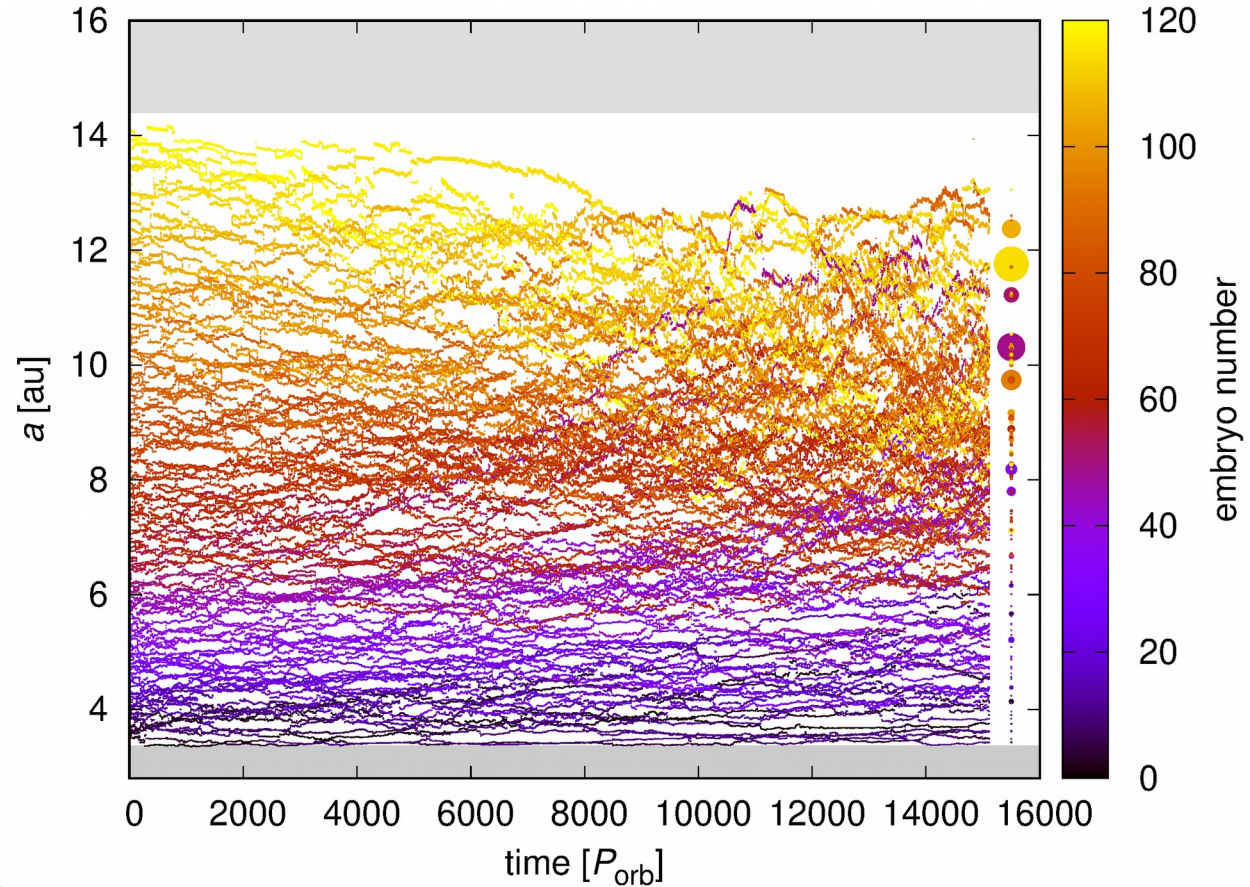
Dynamical tadpole torque

- a vortex behind the protoplanet (Pierens 2015) & its decay → migration cycles



embryos_0.1ME_120

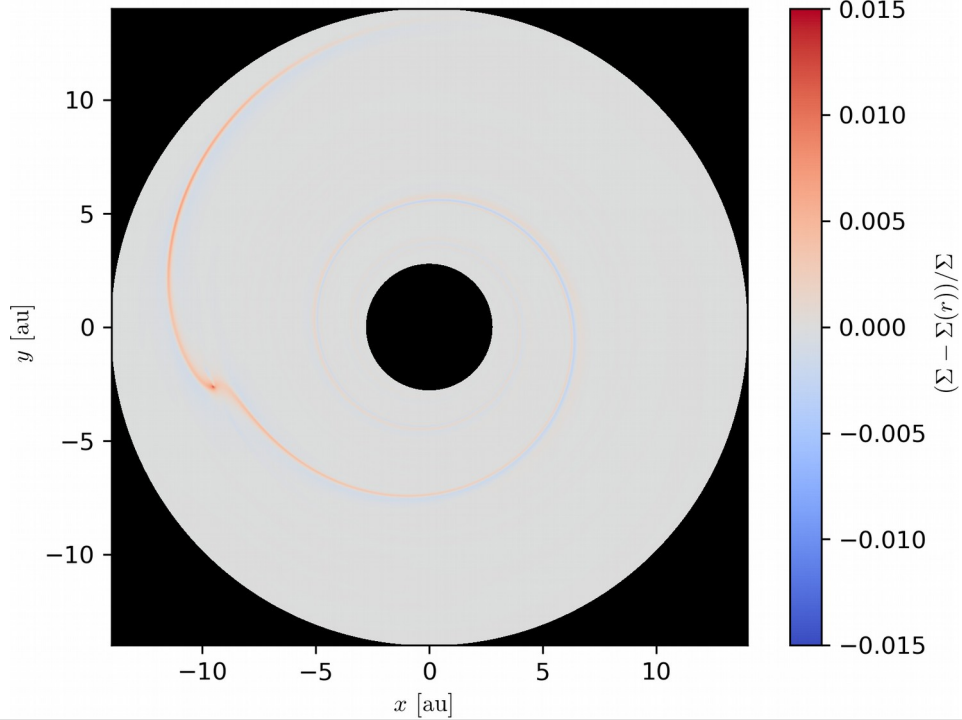
- i.e. Mars-size embryos
- robust runaway growth:
 - (i) merging, (ii) Hill regime of accretion, $dM/dt \sim M_{\text{em}}$,
 - (iii) pebble filtering,
 - (iv) lower inclinations of massive embryos (Levison et al. 2015)
- hot-trail \rightarrow encounters \rightarrow random walk
- BUT also...



Stochastic disk torques

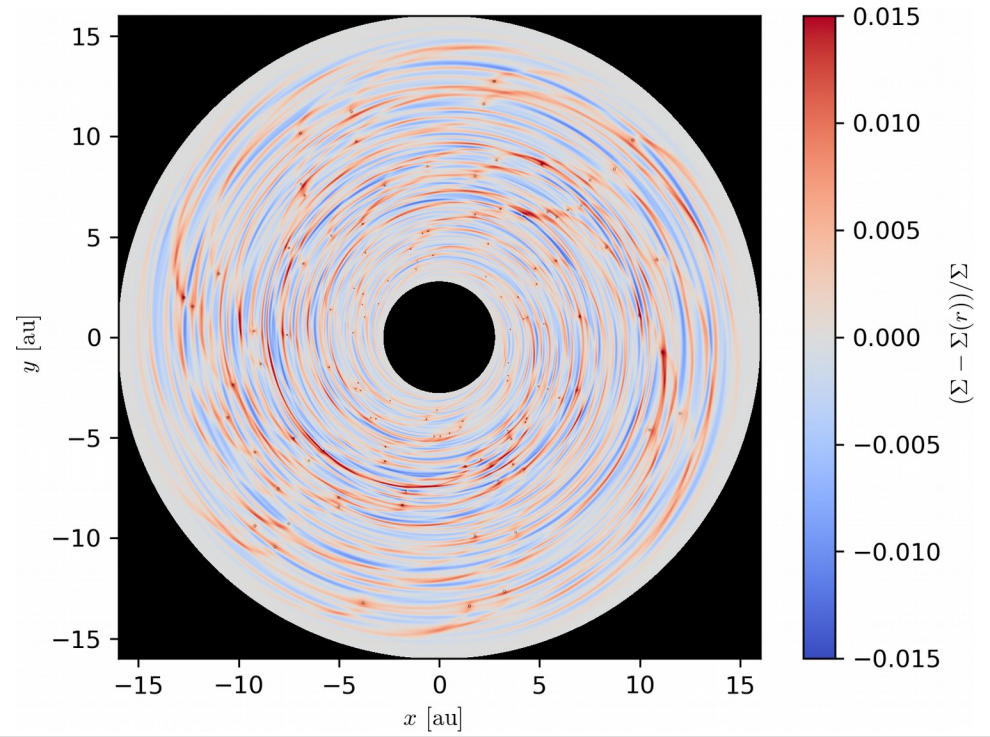
↑ overlapping spiral arms, affected corotation regions, incl. non-linear terms (T^3 , T^4 , Q_{visc})

single-planet torque



vs

stochastic torques



Conclusions

- **3-body encounters** are needed for mergers
- in high- Σ disks **repulsion** events are frequent
- a massive coorbital pair develops a **pebble isolation**
- **stabilisation** & inward migration of c. then occurs
- a **dynamical torque** can arise in the ice-giant zone
- this leads to outward \leftrightarrow inward **migration cycles**
- disk torques for many (>100) low-mass embryos are **stochastic**

- Brož, Chrenko, Nesvorný & Lambrechts (A&A, submitted)
- Fargo-Thorin: <http://sirrah.troja.mff.cuni.cz/~chrenko/>

Inclination vs pebble accretion

- massive embryos \rightarrow low (mean) inclinations, close to pebble disk midplane

