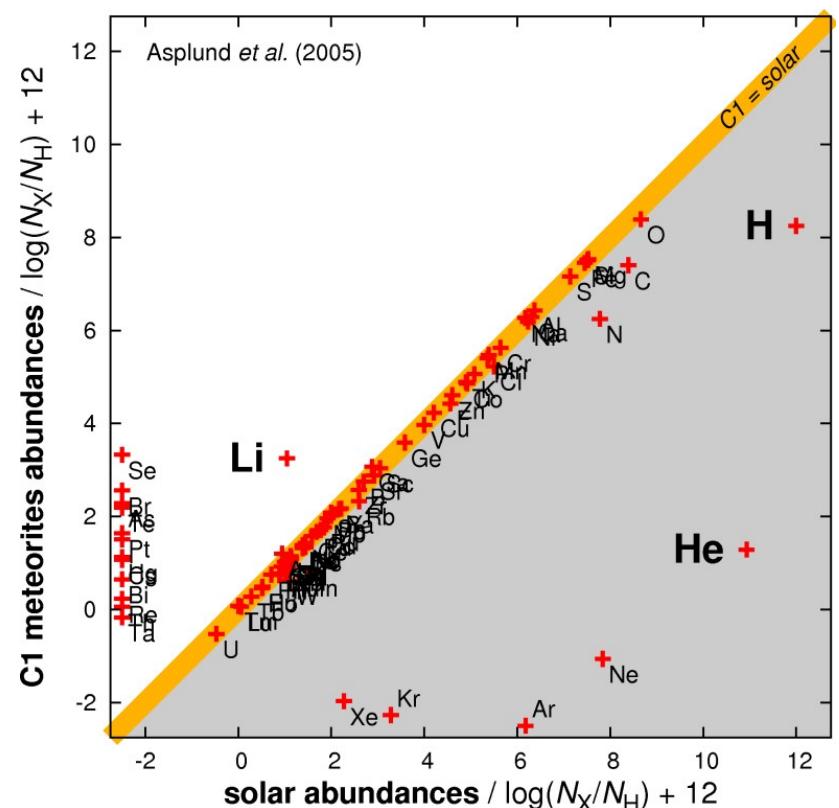


Per asteroides ad astra

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1. general problems
2. observations
3. our dynamical model (1 of)
4. asteroid families
5. future applications



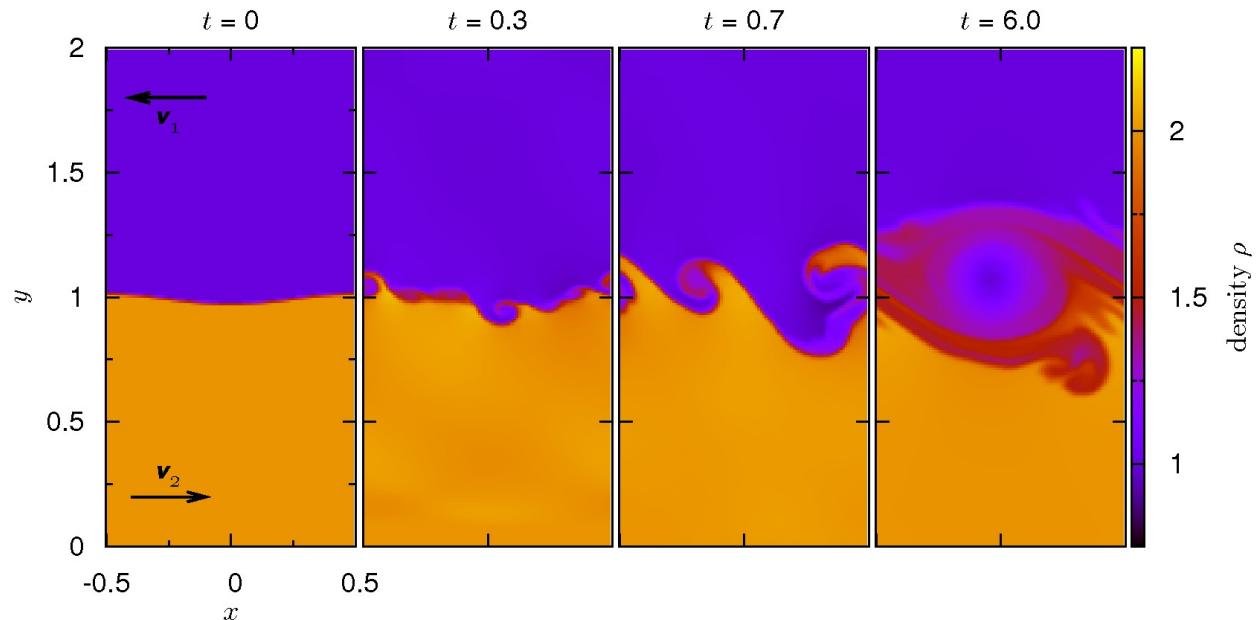
serious

Five problems

- turbulence
- chaos
- irreversibility
- stochasticity
- $t = 0$

\downarrow

an **inverse** problem
(with exceptions)



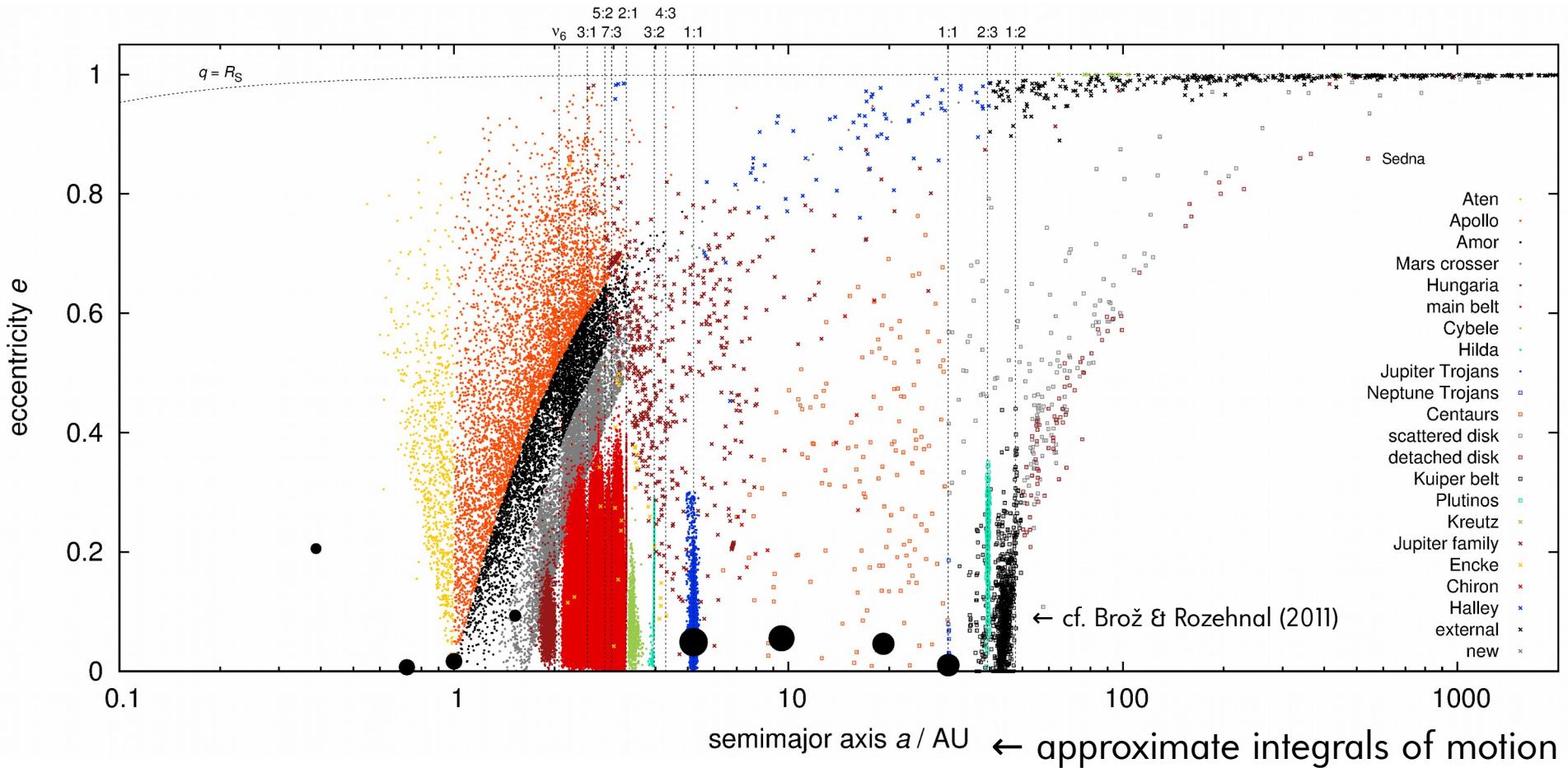
Kelvin–Helmholtz instability
Pluto code (Mignone et al. 2007)

additional instabilities:
Rayleigh–Taylor
magneto-rotational (Flock et al. 2013)
streaming (Johansen et al. 2007)

Observations

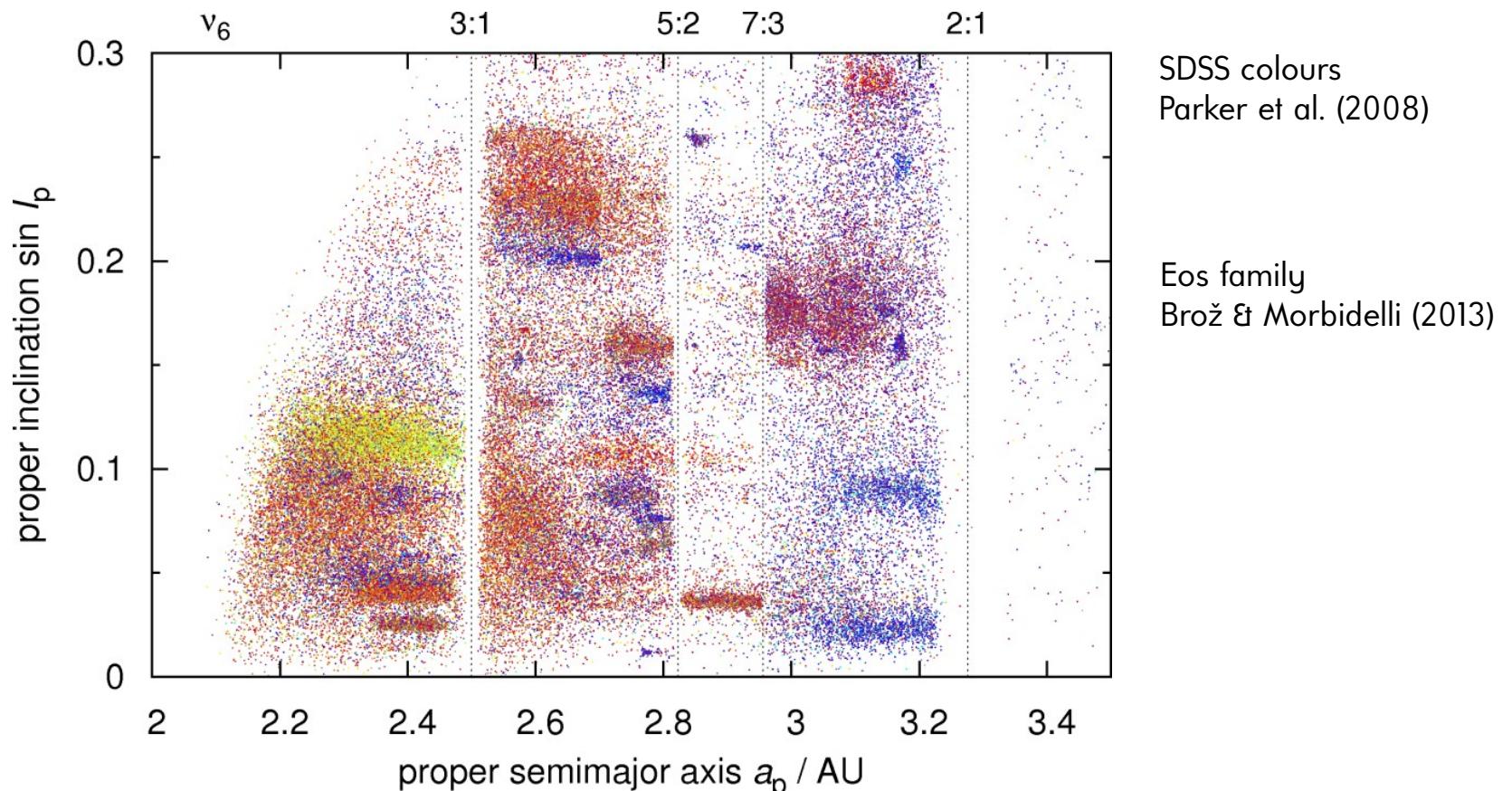
← usually taken @ 4.56 Gyr

- orbital distribution → families everywhere: MB, Hildas, Trojans of J & M, TNOs, irregular moons, ...



Observations (cont.)

- a detail of the Main Asteroid Belt, 124 families in total
(according to our review in AIV, Nesvorný et al. 2015)



a “standard” for families → **N-body model**

We use a symplectic integration scheme (Levison and Duncan 1994), denoted as kick–drift–kick, where the ‘kick’ (actually, a perturbation) is performed as:

$$\dot{\mathbf{r}}_{n+1} = \dot{\mathbf{r}}_n + \ddot{\mathbf{r}} \frac{\Delta t}{2}, \quad (3)$$

and the ‘drift’ corresponds to an analytical solution of the two-body problem (the Sun–asteroid), which involves a numerical solution of the transcendent Kepler equation:

$$M = E - e \sin E, \quad (4)$$

0. drift

$$\mathbf{r}_{n+1} = p(E)\mathbf{r}_n + q(E)\dot{\mathbf{r}}_n, \quad (5)$$

$$\dot{\mathbf{r}}_{n+1} = \dot{p}(E)\mathbf{r}_n + \dot{q}(E)\dot{\mathbf{r}}_n; \quad (6)$$

we account for gravitational perturbations by planets, expressed in the heliocentric frame:

1. kick

$$\ddot{\mathbf{r}}_j = \sum_i \left[-\frac{Gm_i}{r_i^3} \mathbf{r}_i - \frac{Gm_i}{r_{ji}^3} \mathbf{r}_{ji} \right], \quad (7)$$

possibly, the planetary migration, in an analytical way (Malhotra 1995), and also eccentricity damping (Morbidelli et al. 2010):

2. migration

$$\dot{\mathbf{r}}_{n+1} = \dot{\mathbf{r}}_n \left[1 + \frac{\Delta v}{\dot{r}} \frac{\Delta t}{\tau_{\text{mig}}} \exp \left(-\frac{t - t_0}{\tau_{\text{mig}}} \right) \right], \quad (8)$$

N-body model (cont.)

as of Brož et al. (2011)

the Yarkovsky thermal effect (Vokrouhlický 1998, Vokrouhlický and Farinella 1999):

$$f_X(\zeta) + i f_Y(\zeta) = -\frac{8}{3\sqrt{3\pi}} \Phi t'_{1-1}(R'; \zeta), \quad (9)$$

3. IR emission

$$f_Z(\zeta) = -\frac{4}{3} \sqrt{\frac{2}{3\pi}} \Phi t'_{10}(R'; \zeta), \quad (10)$$

$$\Phi \equiv \frac{(1-A)\mathcal{E}_\star \pi R^2}{m_j c_{\text{vac}}}, \quad (11)$$

the YORP effect (Čapek and Vokrouhlický 2004):

$$\dot{\omega} = cf_k(\gamma), \quad (12)$$

4. YORP

$$\dot{\gamma} = \frac{cg_k(\gamma)}{\omega}, \quad (13)$$

$$c \equiv c_{\text{YORP}} \left(\frac{a}{a_0} \right)^{-2} \left(\frac{R}{R_0} \right)^{-2} \left(\frac{\rho}{\rho_0} \right)^{-1}, \quad (14)$$

mass shedding beyond the critical angular frequency (Pravec and Harris 2000):

5. mass shedding

$$\omega_{\text{crit}} = \sqrt{\frac{4}{3}\pi G\rho}, \quad (15)$$

and random collisional reorientations with the time scale (Farinella et al. 1998):

6. collisions

$$\tau_{\text{reor}} = B \left(\frac{\omega}{\omega_0} \right)^{\beta_1} \left(\frac{R}{R_0} \right)^{\beta_2}. \quad (16)$$

A number of unknowns...

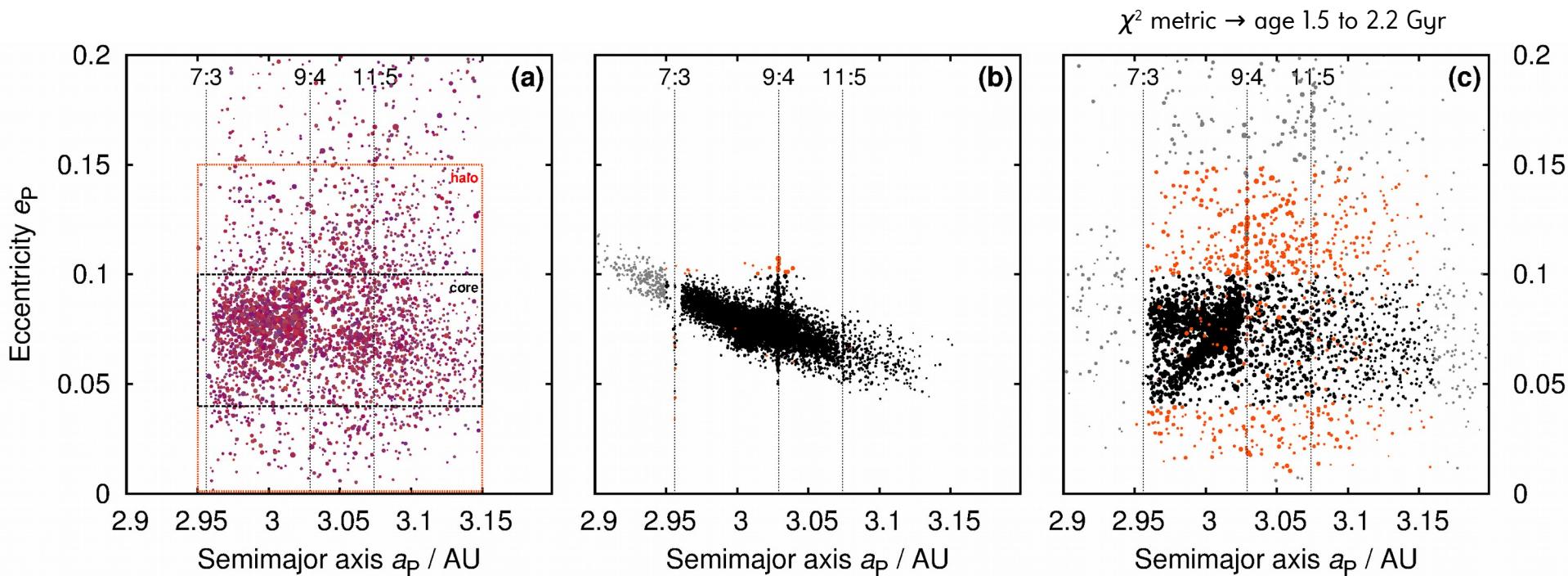
i ... “mass-less” particles, j ... massive bodies

- N_{TP} , \mathbf{r}_i , \mathbf{v}_i , \mathbf{r}_j , \mathbf{v}_j , m_i , m_j , τ_{mig} , Δv , D_i , ρ_i , ρ_{surf} , K , C , A_{Bond} , ϵ , c_{YORP} ,
 λ_i , β_i , ω_i , f_k , g_k , B , β_1 , β_2 , D_0 , D_{PB} , ρ_{PB} , v_{imp} , φ_{imp} , f_{imp} , ω_{imp}
- **32 (!) a-priori unknown ICs and parameters**
- not speaking about Monte-Carlo or SPH models yet...
- time step $\Delta t \rightarrow$ discretisation error \leftarrow usually small(er)
- beware of (formal) uncertainties & (possible) **systematics**

a similar N -body model for multiple stars, e.g. V505 Sgr (Brož et al. 2010), ξ Tau
(Nemráfová et al. in prep.) with χ^2 and simplex to fit minima timings (TTV), radial
velocities RV & speckle-interferometry

Application A: Individual families

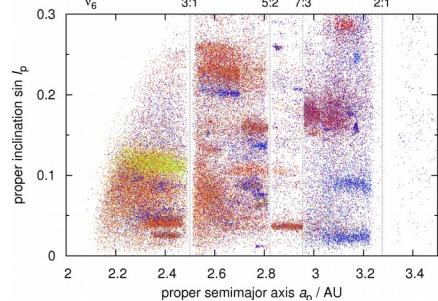
- Eos family (Brož & Morbidelli 2013) → N -body models are *essential* for family identifications!
- core vs halo, K-type taxonomy, distinct from background
- Yarkovsky drift da/dt vs scattering in e , i by resonances



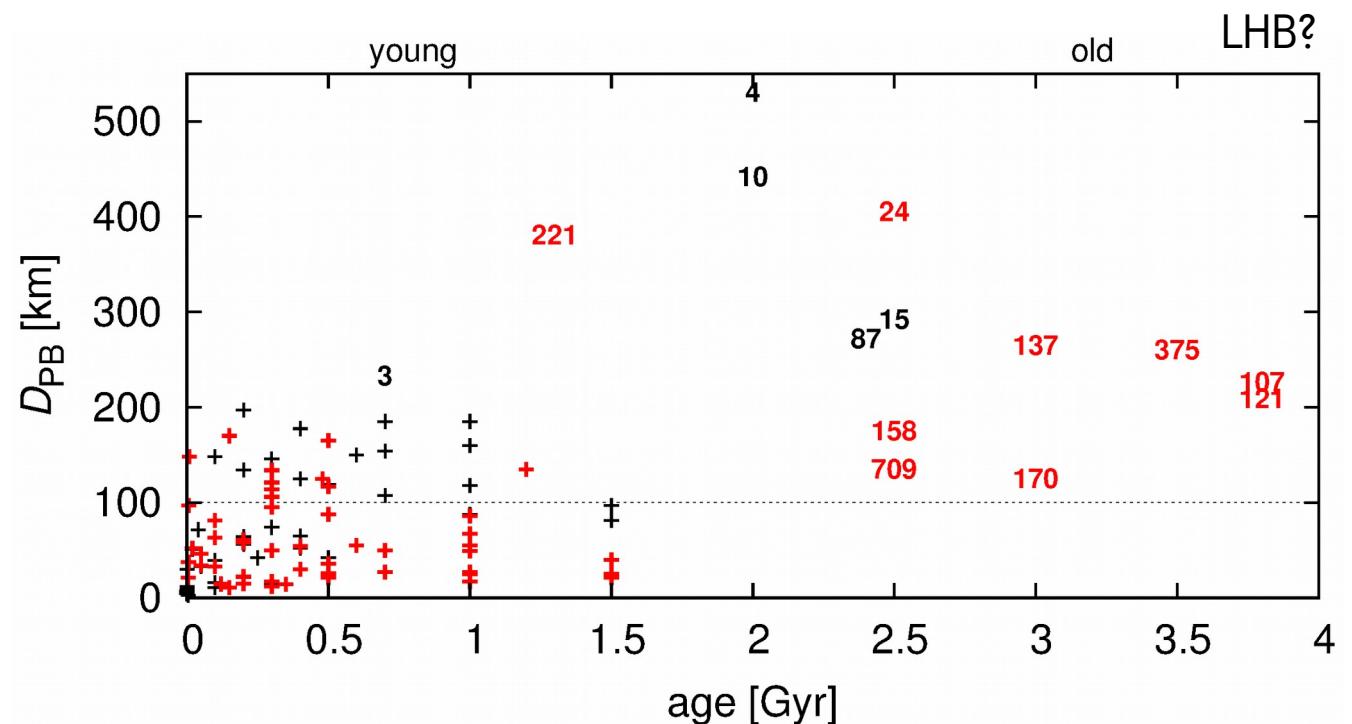
B: Statistics of families

all known, at least

- ages span 4 Gyr (Brož et al. 2013, Bottke et al. 2015) ← OK
- set of **catastrophic disruptions** $D_{\text{PB}} > 100 \text{ km}$ seems complete
- “new” families mostly $D_{\text{PB}} < 100 \text{ km}$, or cratering events

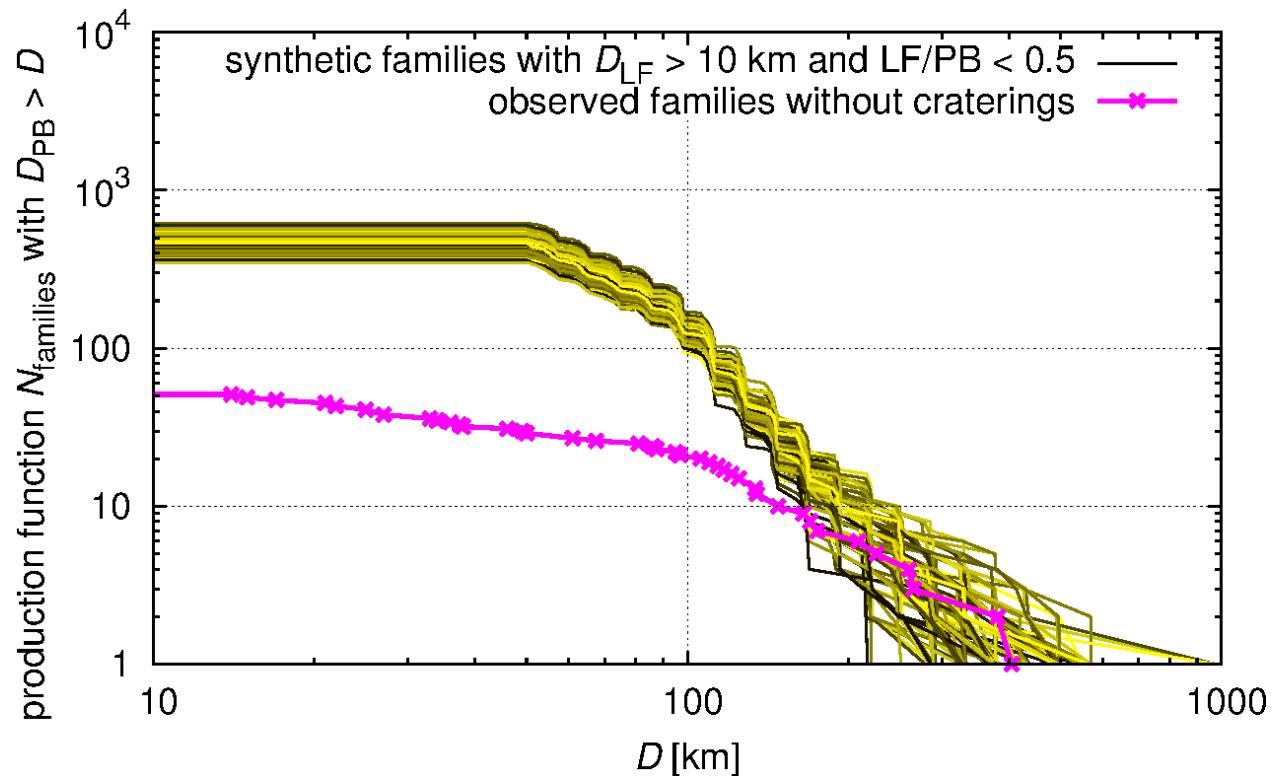
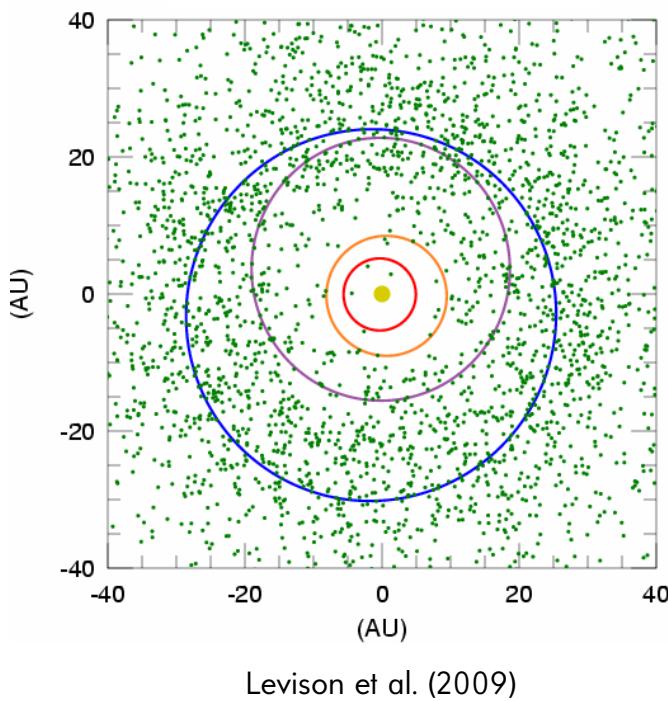


↑
parent-body sizes
estimated by scaling
(Durda et al. 2007)



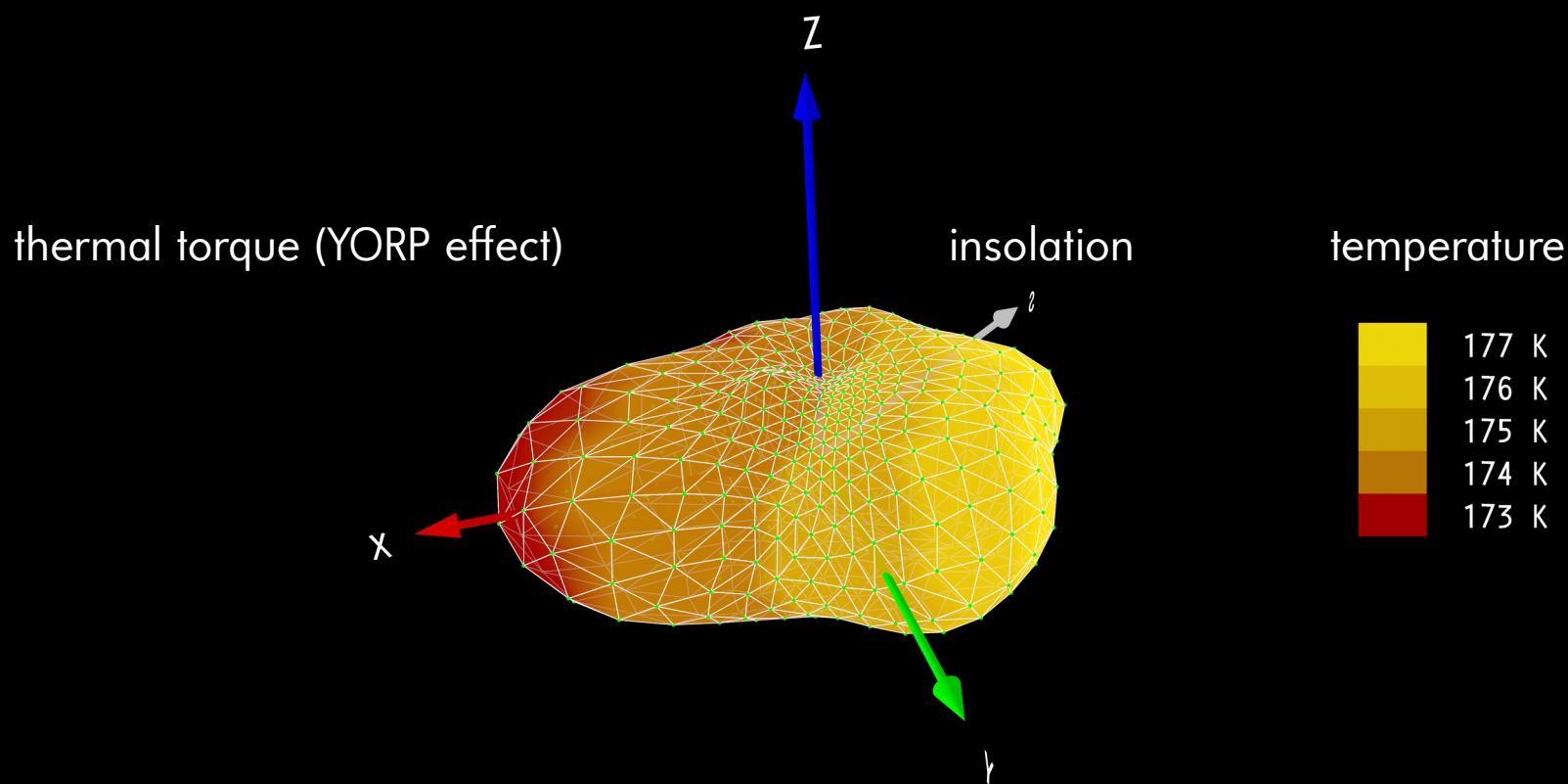
C: Late heavy bombardment of the MB

- no problems producing $D_{\text{PB}} > 200 \text{ km}$ families (Brož et al. 2013)
- *but* 5 times more $D_{\text{PB}} > 100 \text{ km}$ families ← breakups of trans-neptunian comets at low q & secondary collisions



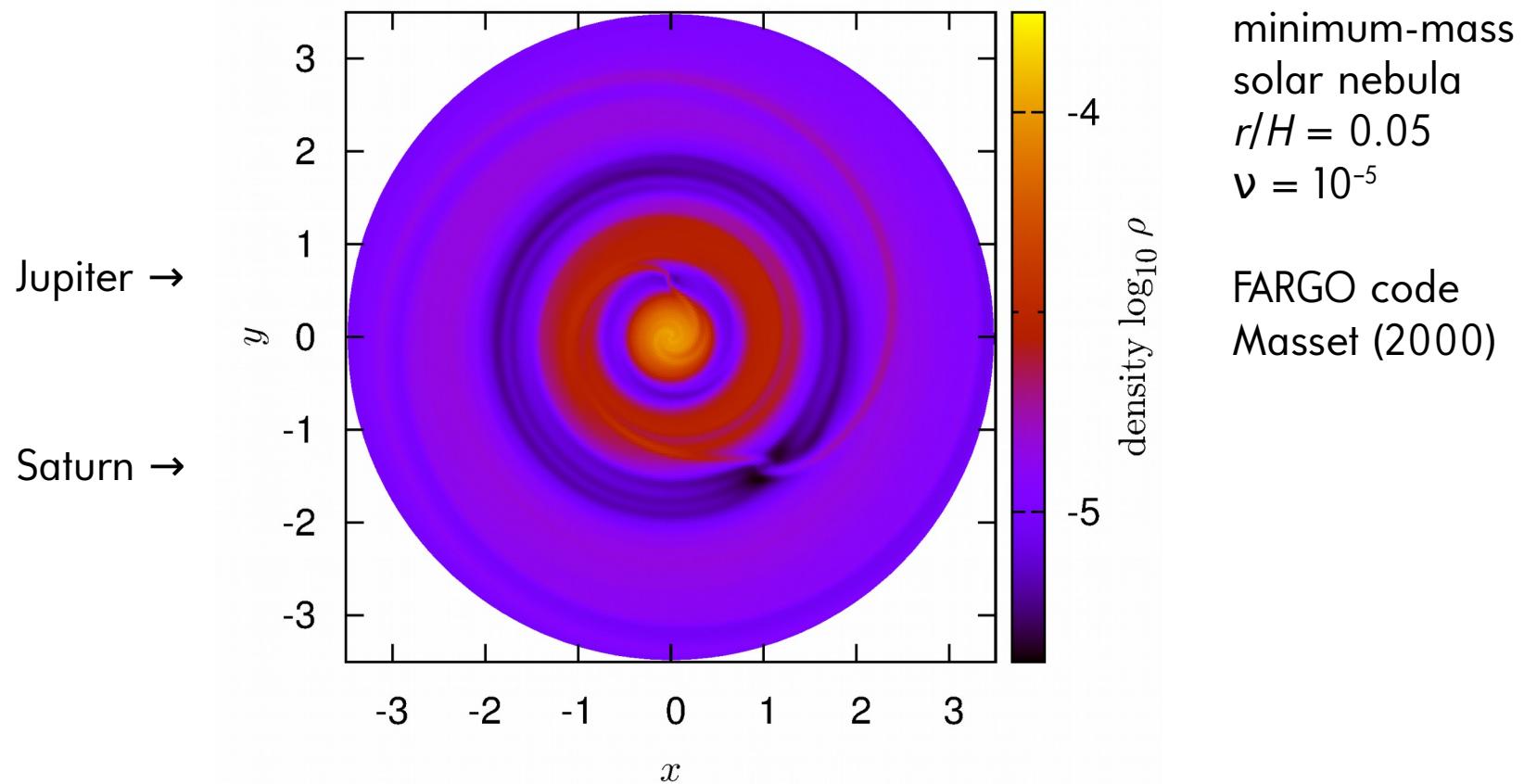
Future applications

- 3-dimensional heat diffusion in meteoroids & boulders (FEM)
- important results for (25147) Itokawa (Ševeček et al. 2015)



Future applications (cont.)

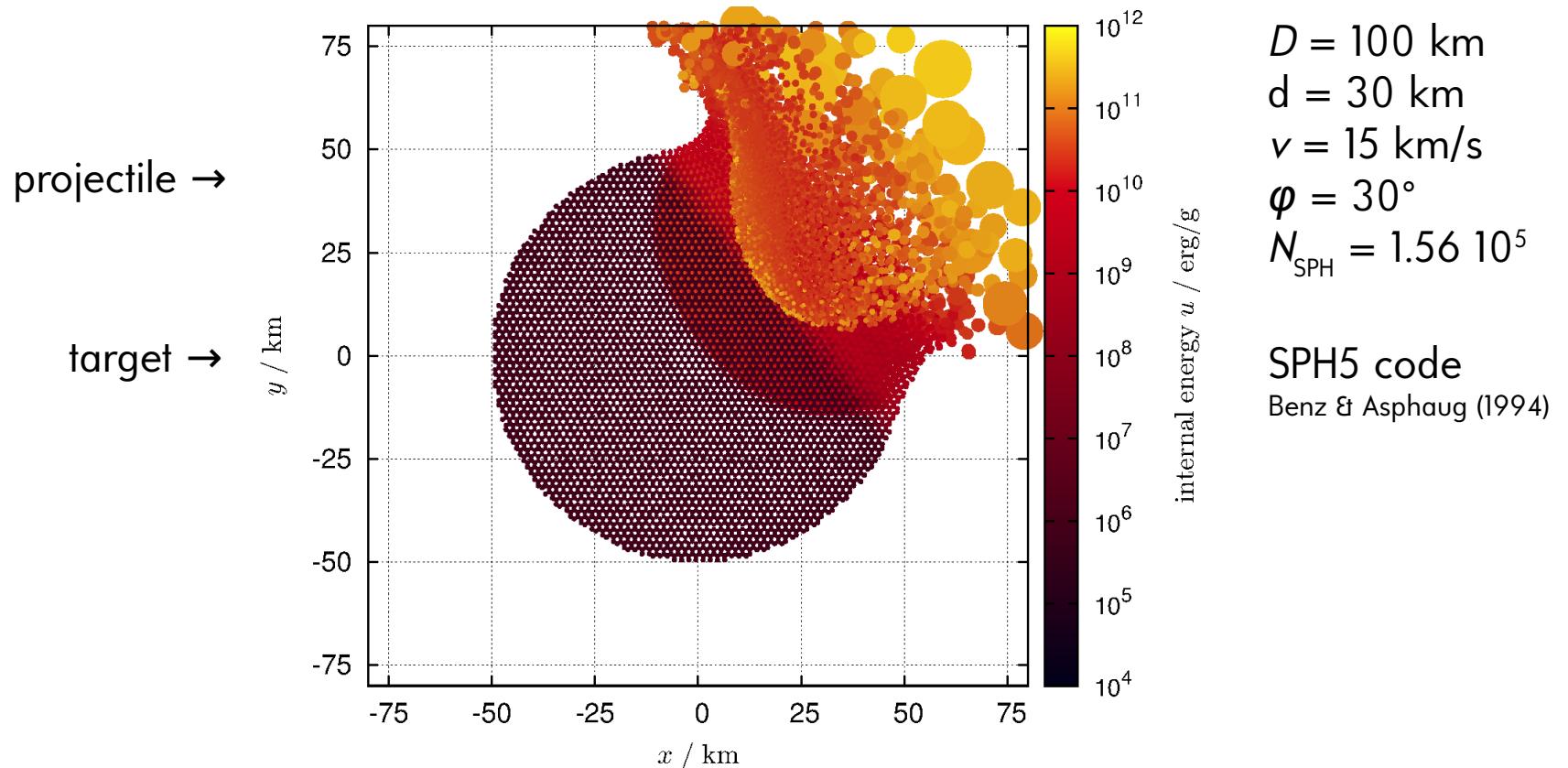
- protoplanetary disks & solid planetesimals vs resonances
- preliminary results in Chrenko & Brož (2015)



Future applications (end)

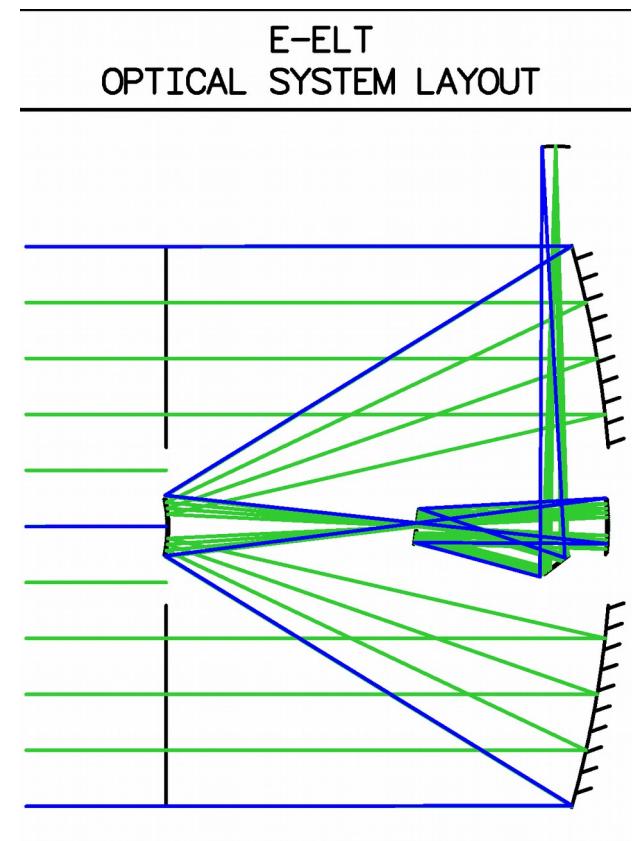
↓ smoothed particle

- SPH simulations of collisions (Rozehnal et al. submitted)
- improve scaling of SPH models ($D_{\text{PB}} >$ and $< 100 \text{ km}$)



Textbooks (in prep.)

- **Hydrodynamics in Astronomy**
protoplanetary disks (FVM),
circumplanetary disks, asteroid
collisions (SPH), cratering, heat
diffusion (FEM), mount elasticity, ...
- **Astronomical Measurements**
statistics, signal to noise, geometrical
optics, diffraction, CCD electronics,
superconductive detectors,
polarimetry, interferometry,
radiotelescopes, particle detectors, ...



Comments of the referees

- asteroids & stars
 - details vs general
 - a convex approximation
-
- paradigm shift (Brož & Rozehnal 2011, Brož & Morbidelli 2013)
 - contradiction vs opportunity (Cibulková et al. 2014)